

Nonisospectral Botie-Pempinelli-Tu hierarchy and its integrable coupling

WANG Hai-feng¹ ZHANG Yu-feng^{2,*}

Abstract. An efficient scheme is applied to generate a nonisospectral Botie-Pempinelli-Tu (BP-T) integrable hierarchy under the case where $\lambda_t = \sum_{j=0}^n k_j(t)\lambda^{-j}$. Based on an expanding higher-dimensional Lie algebra, we obtain a nonisospectral BPT integrable coupling hierarchy. It follows that some new nonisospectral nonlinear systems are obtained by reducing these two nonisospectral BPT hierarchies. Actually, these nonisospectral integrable models that we obtained can enrich the existing integrable models and possibly describe new nonlinear phenomena.

§1 Introduction

In the fields of mathematical physics and theoretical physics, integrable nonlinear systems have attracted wide attention due to the fact that they successfully describe and explain nonlinear phenomena in natural science [1]. Among them, the nonisospectral equations are meaningful and indispensable. A large number of nonisospectral deformations of classic integrable systems are considered [1,2]. Some solitary waves can be described by some nonisospectral equations in a certain type of nonuniform media [3,4]. In [5], the authors presented the nonisospectral linear representations which can help find some geometric properties of some important soliton systems. Additionally, there are many excellent research results of nonisospectral equations related to the inverse scattering transform [6,7]. Therefore, how to generate nonisospectral equations and their extended systems is a crucial research topic [8,9]. A method for generating integrable systems was proposed by Magri [10], which was called the Lax-pair method. Based on it, Tu proposed a method for generating integrable Hamiltonian hierarchies by making use of a trace identity in [11]. Some integrable systems and the corresponding Hamiltonian structures as well as other properties were obtained by using the Tu scheme, such as the work in [12-18]. And Zhang et al deduced an integrable couplings of the well-known BPT hierarchy by constructing a new loop algebra in [19]. There are many different methods for generating isospectral integrable equations. However, as far as we know, fewer works has been done on the nonisospectral integrable equations. In [20,21], Ma proposed a method of constructing its

Received: 2020-04-06. Revised: 2022-04-04.

MR Subject Classification: 37K05, 37K10, 35P30, 37K30.

Keywords: nonisospectral BPT hierarchy, Lie algebra, symmetry, integrable coupling.

Digital Object Identifier(DOI): <https://doi.org/10.1007/s11766-026-4092-0>.

Supported by the National Natural Science Foundation of China (11971475).

*Corresponding author.

corresponding nonisospectral $\lambda_t = \lambda^n (n \geq 0)$ hierarchy of evolution equations by starting from Lax equation $L\psi = \lambda\psi, \lambda_t = \lambda^n$. If $\psi_t = B_n\psi$, then the compatibility condition admits a nonisospectral hierarchy

$$u_t = \sigma_n = \Phi^n g_0, n \geq 0,$$

where the operator Φ is a hereditary operator which satisfying $u_t = K_m(u) = \Phi^m f(\lambda_t = 0)$ [22]. In [23,24], Qiao adopted the Lenard series method to obtain some nonisospectral integrable hierarchies under the case $\lambda_t = \lambda^{m+1}M$.

BPT spatial spectral problem was first introduced by Boiti, Pempinelli and Tu in [25], where they considered a particular case of the more general spectral problem proposed by Boiti and Tu in [26] and they showed that the soliton equations of the related hierarchy are Hamiltonian systems with commuting flows on a symplectic Kahler manifold. In this article, an efficient scheme is applied to generate a nonisospectral BPT integrable hierarchy under the case where $\lambda_t = \sum_{j=0}^n k_j(t)\lambda^{-j}$ [27]. One can find that the case is a generalized expression for the case $\lambda_t = \lambda^n$. Based on it, we deduce an expanding nonisospectral integrable hierarchies by using the knowledge of integrable coupling [28]. In [29,30], Ma presented the integrable couplings can be derived based on the corresponding non-semisimple Lie algebra \bar{g} which has the following triangular block matrix form

$$M(A, B) = \begin{bmatrix} A & 0 \\ B & A \end{bmatrix}, \quad (1)$$

where A and B are two arbitrary square matrices of the same order. By reducing the nonisospectral BPT integrable hierarchy and its integrable coupling hierarchy, we obtain some new nonisospectral nonlinear systems. Ma proposed an effective method to calculate symmetry algebras and successfully applied it to coupled KdV systems [31]. Then, the types of algebraic structures of the spaces of the isospectral Lax operators were established, and it means that the theoretical basis of the Lax operator method has been essentially formed [32]. Under obtaining the expanding nonisospectral integrable systems, some properties including symmetry algebras, Bäcklund transformations [34], analytic solutions [35], and so on, could be studied [36]. The research object of this paper is the generation of the nonisospectral integrable hierarchy and its extended hierarchy, which is different from the related research on the isospectral hierarchy. As far as we know, the nonisospectral BPT hierarchy and its extended integrable hierarchy obtained in this paper are new and different from the existing literature.

§2 A nonisospectral BPT hierarchy

A basis of the loop algebras \bar{A}_1 is given by

$$\begin{cases} h(n) = \frac{1}{2} \begin{pmatrix} \lambda^n & 0 \\ 0 & -\lambda^n \end{pmatrix}, & e(n) = \frac{1}{2} \begin{pmatrix} 0 & \lambda^n \\ \lambda^n & 0 \end{pmatrix}, & f(n) = \frac{1}{2} \begin{pmatrix} 0 & \lambda^n \\ -\lambda^n & 0 \end{pmatrix}, \\ \deg h(n) = \deg e(n) = \deg f(n) = n, \end{cases} \quad (2)$$

equipped with

$$[h(n), e(m)] = f(m+n), \quad [h(n), f(m)] = e(m+n), \quad [e(n), f(m)] = -h(m+n).$$

Based on the loop algebra $\overline{A_1}$, we consider the corresponding nonisospectral problems

$$\begin{cases} \psi_x = U\psi, & U = 2h(1) + qe(0) + rf(-1) + sh(-1) = \begin{pmatrix} \lambda + \frac{1}{2}s\lambda^{-1} & \frac{1}{2}q + \frac{1}{2}r\lambda^{-1} \\ \frac{1}{2}q - \frac{1}{2}r\lambda^{-1} & -\lambda - \frac{1}{2}s\lambda^{-1} \end{pmatrix}, \\ \psi_t = V\psi, & V = \sum_{m \geq 0} (a_m h(-m) + b_m e(-m) + c_m f(-m)), \\ \lambda_t = \sum_{m \geq 0} k_m(t) \lambda^{1-m}. \end{cases} \quad (3)$$

In order to deduce the nonisospectral integrable hierarchy, one need solve the nonisospectral zero curvature equation

$$\frac{\partial U}{\partial u} u_t + \frac{\partial U}{\partial \lambda} \lambda_t^{(n)} - V_x^{(n)} + [U, V^{(n)}] = 0. \quad (4)$$

From (3), we have

$$\begin{aligned} \frac{\partial U}{\partial \lambda} \lambda_t &= \begin{pmatrix} 1 - \frac{1}{2}s\lambda^{-2} & -\frac{1}{2}r\lambda^{-2} \\ \frac{1}{2}r\lambda^{-2} & -1 + \frac{1}{2}s\lambda^{-2} \end{pmatrix} \sum_{m \geq 0} k_m(t) \lambda^{1-m} \\ &= \sum_{m \geq 0} k_m(t) [2h(1-m) - rf(-1-m) - sh(-1-m)]. \end{aligned}$$

By solving the stationary zero curvature equation

$$V_x = \frac{\partial U}{\partial \lambda} \lambda_t + [U, V], \quad (5)$$

we obtain the recursion relations

$$\begin{cases} a_{nx} = 2k_{n+1}(t) - sk_{n-1}(t) + rb_{n-1} - qc_n, \\ b_{nx} = 2c_{n+1} + sc_{n-1} - ra_{n-1}, \\ c_{nx} = -rk_{n-1}(t) + 2b_{n+1} + sb_{n-1} - qa_n. \end{cases} \quad (6)$$

Take $a_0 = 2\beta$, $b_0 = c_0 = 0$, and then one has

$$\begin{aligned} b_1 &= \beta q, \quad c_2 = \frac{\beta q_x}{2} + \beta r, \quad a_2 = -\frac{\beta q^2}{4} + 2k_3(t)x - sk_1(t)x + \beta_1, \dots, \\ a_{2m+1} &= b_{2m} = c_{2m+1} = k_{2m}(t) = 0. \end{aligned}$$

Note that

$$\begin{aligned} V_+^{(n)} &= \sum_{m=0}^n (a_m h(-m) + b_m e(-m) + c_m f(-m)) \lambda^n = \lambda^n V - V_-^{(n)}, \\ \lambda_{t,+}^{(n)} &= \lambda^n \lambda_t - \lambda_{t,-}^{(n)} = \sum_{m=0}^n k_m(t) \lambda^{n-m+1}, \end{aligned}$$

then (5) can be broken down into

$$-V_{+,x}^{(n)} + \frac{\partial U}{\partial \lambda} \lambda_{t,+}^{(n)} + [U, V_+^{(n)}] = V_{-,x}^{(n)} - \frac{\partial U}{\partial \lambda} \lambda_{t,-}^{(n)} - [U, V_-^{(n)}]. \quad (7)$$

It follows that

$$\begin{aligned} \deg V_{+,x}^{(n)} &=: (0, 0, 0), \quad \deg \frac{\partial U}{\partial \lambda} \lambda_{t,+}^{(n)} =: (1, -1, -1), \quad \deg([U, V_+^{(n)}]) =: (1, 0, -1, -1; 0, 0, 0), \\ \deg V_{-,x}^{(n)} &=: (-1, -1, -1), \quad \deg \frac{\partial U}{\partial \lambda} \lambda_{t,-}^{(n)} =: (0, -2, -2), \quad \deg([U, V_-^{(n)}]) =: (1, 0, -1, -1; -1, -1, -1). \end{aligned}$$

We find that the gradation of the left-hand side of (7) is more than -1 , while the right-hand side is less than 0 . Therefore, take the gradations 0 and -1 in (7), one has

$$\begin{aligned}
 & -V_{+,x}^{(n)} + \frac{\partial U}{\partial \lambda} \lambda_{t,+}^{(n)} + [U, V_+^{(n)}] = -2c_{n+1}e(0) + (-b_{n+1,x} + 2c_{n+2})e(-1) + 2b_{n+1}f(0) \\
 & + (c_{n+1,x} - 2b_{n+2} + qa_{n+1})f(-1) + 2k_{n+1}(t)h(0) + (a_{n+1,x} + qc_{n+1} - 2k_{n+2}(t))e(-1).
 \end{aligned}$$

Choose $n = 2m - 1$, one can find $\Delta_n = 0$ so that $V^{(n)} = V_+^{(n)}$.

Theorem 1. By means of (4) and (6), we obtain the nonisospectral BPT hierarchy as follows:

$$\begin{aligned}
 u_t = \begin{pmatrix} q \\ r \\ s \end{pmatrix}_t &= \begin{pmatrix} 2c_{2m} \\ -sb_{2m-1} + rk_{2m-1}(t) \\ sk_{2m-1}(t) - rb_{2m-1} \end{pmatrix} = J_1 \begin{pmatrix} b_{2m+1} \\ -c_{2m} \\ a_{2m} \end{pmatrix} + \begin{pmatrix} 0 \\ 0 \\ 2k_{2m+1}(t) \end{pmatrix} \\
 &= J_1 L_1^m \begin{pmatrix} \beta q \\ 0 \\ 2\beta \end{pmatrix} + J_1 \sum_{i=0}^{m-1} (L_1^i K_{2m+1-2i}(t)xQ_1) + J_1 \sum_{i=0}^{m-1} (L_1^i K_{2m-1-2i}(t)Q_2) \\
 &\quad + J_1 \sum_{i=0}^{m-1} (L_1^i \alpha_{2m-2i}R_1) + \begin{pmatrix} 0 \\ 0 \\ 2k_{2m+1}(t) \end{pmatrix} \tag{8} \\
 &= \Phi^m J_1 \begin{pmatrix} \beta q \\ 0 \\ 2\beta \end{pmatrix} + \sum_{i=0}^{m-1} (\Phi^i J_1 K_{2m+1-2i}(t)xQ_1) + \sum_{i=0}^{m-1} (\Phi^i J_1 K_{2m-1-2i}(t)Q_2) \\
 &\quad + \sum_{i=0}^{m-1} (\Phi^i J_1 \alpha_{2m-2i}R_1) + \begin{pmatrix} 0 \\ 0 \\ 2k_{2m+1}(t) \end{pmatrix},
 \end{aligned}$$

where

$$\begin{aligned}
 J_1 &= \begin{pmatrix} 0 & -2 & 0 \\ 2 & \partial & -q \\ 0 & q & -\partial \end{pmatrix}, \quad Q_1 = \begin{pmatrix} q \\ 0 \\ 2 \end{pmatrix}, \quad Q_2 = \frac{1}{2} \begin{pmatrix} -qsx + r \\ 0 \\ -2sx \end{pmatrix}, \quad R_1 = \frac{1}{4} \begin{pmatrix} q \\ 0 \\ 1 \end{pmatrix}, \\
 L_1 &= \frac{1}{4} \begin{pmatrix} \partial^2 - 2s - q^2 + 2q\partial^{-1}r + q\partial^{-1}q_x & \partial s - q\partial^{-1}(qs) & \partial r - q\partial^{-1}(qr) \\ & -2\partial & -2s & -2r \\ -2q + \partial^{-1}(2q_x + 4r) & -2\partial^{-1}(qs) & -2\partial^{-1}(qr) \end{pmatrix},
 \end{aligned}$$

with

$$\Phi = J_1 L_1 J_1^{-1} = \frac{1}{4} \begin{pmatrix} -q_x \partial^{-1}q - q^2 + \partial^2 - 2s - 2r\partial^{-1}q & 0 & -2q_x \partial^{-1} - 2q - 2r\partial^{-1} \\ & qs\partial^{-1}q - s\partial & -2s & 2qs\partial^{-1} \\ & qr\partial^{-1}q - r\partial & -2r & 2qr\partial^{-1} \end{pmatrix}. \tag{9}$$

To the nonisospectral BPT hierarchy (8), we show the first two nonlinear systems as follows:

$$\begin{cases} q_t = \beta q_x + 2\beta r, \\ r_t = -s\beta q + rk_1(t), \\ s_t = -r\beta q + sk_1(t), \end{cases} \tag{10}$$

and

$$\begin{cases} q_t = \frac{\beta}{4}(q_{xxx} + 2r_{xx} - 4q_x s - 2qs_x - \frac{3}{2}q^2 q_x + 4rs - q^2 r) + \frac{\beta_1}{2}(q_x + 2r) \\ \quad + \frac{k_1(t)}{2}(r_x - qs - xq_x s - xqs_x - 2xrs) + k_3(t)(q + xq_x + 2xr), \\ r_t = -\frac{\beta}{4}(q_{xx} s + 2r_x s - 2qs^2 - \frac{1}{2}q^3 s) - \frac{\beta_1}{2}qs - \frac{k_1(t)}{2}(rs - xqs^2) + k_3(t)(xqs + r), \\ s_t = -\frac{\beta}{4}(q_{xx} s + 2r_x r - 2qsr - \frac{1}{2}q^3 r) - \frac{\beta_1}{2}qr - \frac{k_1(t)}{2}(r^2 - xqsr) + k_3(t)(xqr + s). \end{cases} \quad (11)$$

§3 An expanding nonisospectral BPT hierarchy

The Lie algebras $A_1 = \text{span}\{h, e, f\}$ can be expanded into the Lie algebras $A_{11} = \text{span}\{f_i\}_{i=1}^5$, where

$$f_1 = \frac{1}{2} \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad f_2 = \frac{1}{2} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad f_3 = \frac{1}{2} \begin{pmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix},$$

$$f_4 = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 1 \\ 0 & 0 & 0 \end{pmatrix}, \quad f_5 = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & -1 \\ 0 & 0 & 0 \end{pmatrix},$$

$$[f_1, f_2] = f_3, \quad [f_1, f_3] = f_2, \quad [f_2, f_3] = -f_1, \quad [f_1, f_4] = \frac{1}{2}f_5, \quad [f_1, f_5] = \frac{1}{2}f_4,$$

$$[f_2, f_4] = \frac{1}{2}f_4, \quad [f_2, f_5] = -\frac{1}{2}f_5, \quad [f_3, f_4] = \frac{1}{2}f_5, \quad [f_3, f_5] = -\frac{1}{2}f_4, \quad [f_4, f_5] = 0.$$

To the Lie algebra A_{11} , the following loop algebra can be introduced

$$\tilde{A}_{11} = \text{span}\{f_i(n)\}_{i=1}^5, \quad f_i(n) = f_i \lambda^n, \quad [f_i(n), f_j(m)] = [f_i, f_j] \lambda^{m+n}, \quad i, j = 1, \dots, 5, \quad m, n \in \mathbb{Z}.$$

From \tilde{A}_{11} , we consider a nonisospectral problem

$$\begin{cases} \psi_x = M\psi, \quad M = 2f_1(1) + u_1 f_2(0) + u_2 f_3(-1) + u_3 f_1(-1) + u_4 f_4(-1) + u_5 f_5(0), \\ \psi_t = N\psi, \quad N = \sum_{m \geq 0} (a_m f_1(-m) + b_m f_2(-m) + c_m f_3(-m) + d_m f_4(-m) + e_m f_5(-m)), \\ \lambda_t = \sum_{m \geq 0} k_m(t) \lambda^{-m}, \end{cases} \quad (12)$$

then the stationary zero curvature equation admits that

$$\begin{cases} a_{nx} = 2k_{n+1}(t) - u_3 k_{n-1}(t) + u_2 b_{n-1} - u_1 c_n, \\ b_{nx} = 2c_{n+1} - u_2 a_{n-1} + u_3 c_{n-1}, \\ c_{nx} = -u_2 k_{n-1}(t) + 2b_{n+1} - u_1 a_n + u_3 b_{n-1}, \\ d_{nx} = -u_4 k_{n-1}(t) + e_{n+1} + \frac{1}{2}u_1 d_n - \frac{1}{2}u_2 e_{n-1} + \frac{1}{2}u_3 e_{n-1} - \frac{1}{2}u_4 b_{n-1} - \frac{1}{2}u_5 a_n + \frac{1}{2}u_5 c_n, \\ e_{nx} = d_{n+1} - \frac{1}{2}u_1 e_n + \frac{1}{2}u_2 d_{n-1} + \frac{1}{2}u_3 d_{n-1} - \frac{1}{2}u_4 a_{n-1} - \frac{1}{2}u_4 c_{n-1} + \frac{1}{2}u_5 b_n. \end{cases} \quad (13)$$

Take $a_0 = 2\beta$, $b_0 = c_0 = d_0 = e_0 = 0$, then

$$b_1 = \beta u_1, \quad e_1 = \beta u_5, \quad c_2 = \frac{\beta u_1 x}{2} + \beta u_2, \quad a_2 = -\frac{\beta u_1^2}{4} + 2k_3(t)x - u_3 k_1(t)x + \beta_1,$$

$$d_2 = \beta u_5 x + 2\beta u_4, \quad \dots, \quad a_{2m+1} = b_{2m} = c_{2m+1} = e_{2m} = d_{2m+1} = k_{2m}(t) = 0.$$

When $n = 2m - 1$, $N^{(n)} = N_+^{(n)} = \sum_{m=0}^n (a_m f_1(n-m) + b_m f_2(n-m) + c_m f_3(n-m) + d_m f_4(n-m) + e_m f_5(n-m))$.

Theorem 2. From (12) and (13), we obtain the expanding nonisospectral BPT hierarchy

$$\begin{aligned}
 u_t = \begin{pmatrix} u_1 \\ u_2 \\ u_3 \\ u_4 \\ u_5 \end{pmatrix}_t &= \begin{pmatrix} 2c_{2m} \\ -u_3b_{2m-1} + u_2k_{2m-1}(t) \\ -u_2b_{2m-1} + u_3k_{2m-1}(t) \\ \frac{1}{2}(u_2 - u_3)e_{2m-1} + \frac{1}{2}u_4b_{2m-1} + u_4k_{2m-1}(t) \\ d_{2m} \end{pmatrix} = J_2 \begin{pmatrix} b_{2m+1} \\ -c_{2m} \\ a_{2m} \\ e_{2m+1} \\ d_{2m} \end{pmatrix} + \begin{pmatrix} 0 \\ 0 \\ 2k_{2m+1}(t) \\ 0 \\ 0 \end{pmatrix} \\
 &= J_2 L_2^m \begin{pmatrix} \beta u_1 \\ 0 \\ 2\beta \\ \beta u_5 \\ 0 \end{pmatrix} + J_2 \sum_{i=0}^{m-1} (L_2^i K_{2m+1-2i}(t)xQ_3) + J_2 \sum_{i=0}^{m-1} (L_2^i K_{2m-1-2i}(t)Q_4) \\
 &\quad + J_2 \sum_{i=0}^{m-1} (L_2^i \alpha_{2m-2i} R_2) + \begin{pmatrix} 0 \\ 0 \\ 2k_{2m+1}(t) \\ 0 \\ 0 \end{pmatrix} \tag{14} \\
 &= \Phi_2^m J_2 \begin{pmatrix} \beta u_1 \\ 0 \\ 2\beta \\ \beta u_5 \\ 0 \end{pmatrix} + \sum_{i=0}^{m-1} (\Phi_2^i J_2 K_{2m+1-2i}(t)xQ_3) + \sum_{i=0}^{m-1} (\Phi_2^i J_2 K_{2m-1-2i}(t)Q_4) \\
 &\quad + \sum_{i=0}^{m-1} (\Phi_2^i J_2 \alpha_{2m-2i} R_2) + \begin{pmatrix} 0 \\ 0 \\ 2k_{2m+1}(t) \\ 0 \\ 0 \end{pmatrix},
 \end{aligned}$$

where

$$J_2 = \begin{pmatrix} 0 & -2 & 0 & 0 & 0 \\ 2 & \partial & -u_1 & 0 & 0 \\ 0 & u_1 & -\partial & 0 & 0 \\ 0 & -\frac{u_5}{2} & -\frac{u_5}{2} & 1 & \frac{u_1}{2} - \partial \\ 0 & 0 & 0 & 0 & 1 \end{pmatrix}, \quad Q_3 = \begin{pmatrix} u_1 \\ 0 \\ 2 \\ u_5 \\ 0 \end{pmatrix}, \quad Q_4 = \frac{1}{2} \begin{pmatrix} -u_1 u_3 x + u_2 \\ 0 \\ -2u_3 x \\ -u_3 u_5 x + 2u_4 \\ 0 \end{pmatrix}, \quad R_2 = \frac{1}{4} \begin{pmatrix} 2u_1 \\ 0 \\ 2 \\ u_5 \\ 0 \end{pmatrix},$$

$$\Phi_2 = J_2 L_2 J_2^{-1},$$

$$L_2 = \frac{1}{4} \begin{pmatrix} A & \partial u_3 - u_1 \partial^{-1}(u_1 u_3) & \partial u_2 - u_1 \partial^{-1}(u_1 u_2) & 0 & 0 \\ -2\partial & -2u_3 & -2u_2 & 0 & 0 \\ -2u_1 + 2\partial^{-1}(u_1 x + 2u_2) & -2\partial^{-1}(u_1 u_3) & -2\partial^{-1}(u_1 u_2) & 0 & 0 \\ B & C & D & E & F \\ -2u_5 & -2u_4 & 2u_4 & 4\partial + 2u_1 & -2(u_2 + u_3) \end{pmatrix},$$

with

$$\begin{cases} A = \partial^2 - 2u_3 - u_1^2 + 2u_1 \partial^{-1} u_2 + u_1 \partial^{-1} u_{1x}, & B = -2\partial u_5 + 2u_4 - u_5 \partial + u_5 \partial^{-1} u_{1x} + 2u_5 \partial^{-1} u_2, \\ C = -2\partial u_4 + u_1 u_4 + u_3 u_5 + u_5 \partial^{-1} u_1 u_3, & D = 2\partial u_4 - u_1 u_4 - u_2 u_5 - u_5 \partial^{-1} u_1 u_2, \\ E = 4\partial^2 + 2\partial u_1 - 2u_1 \partial - u_1^2 + 2u_2 - 2u_3, & F = -2\partial(u_2 + u_3) + u_1 u_2 + u_1 u_3. \end{cases}$$

To the extended nonisospectral BPT hierarchy (14), the first nonlinear system is

$$\begin{cases} u_{1t} = \beta u_{1x} + 2\beta u_2, \\ u_{2t} = -\beta u_1 u_3 + k_1(t)u_2, \\ u_{3t} = -\beta u_1 u_2 + k_1(t)u_3, \\ u_{4t} = \frac{\beta}{2}(u_2 u_5 + u_1 u_4 - u_3 u_5) + k_1(t)u_4, \\ u_{5t} = \beta u_{5x} + 2\beta u_4, \end{cases} \quad (15)$$

which is an extended system of the system (10). Obviously, the hierarchy (14) is an integrable coupling of the nonisospectral BPT hierarchy (8).

§4 Conclusions and discussions

This paper introduced the nonisospectral problems (3) based on the loop algebras \overline{A}_1 , and thus deduced the nonisospectral BPT hierarchy (8). The nonisospectral problem (12) was introduced by using the extended Lie algebras A_{11} , and then obtained the expanded nonisospectral BPT hierarchy (14). By reducing these two nonisospectral BPT hierarchies, some new nonisospectral nonlinear systems were obtained. Actually, these nonisospectral integrable model that we obtained can enrich the existing integrable models and possibly describe new nonlinear phenomena [37-44].

Declarations

Conflict of interest The authors declare no conflict of interest.

References

- [1] D Levi, O Ragnisco. *Non-isospectral deformations and Darboux transformations for the third-order spectral problem*, Inverse Probl, 1988, 4(3): 815.
- [2] P R Gorda, A Pickering. *Nonisospectral scattering problems: A key to integrable hierarchies*, J Math Phys, 1999, 40(11): 5749-5786.
- [3] F Calogero. *Bäcklund transformations and functional relation for solutions of nonlinear partial differential equations solvable via the inverse scattering method*, Lett Nuovo Cimento, 1975, 14(15): 537-543.
- [4] F Calogero, A Degasperis. *Solution by the spectral-transform method of a nonlinear evolution equation including as a special case the cylindrical KdV equation*, Lett Nuovo Cimento, 1978, 123: 150-154.
- [5] C Rogers, W K Schief. *Bäcklund and Darboux Transformations: Geometry and Modern Applications in Soliton Theory*, Cambridge: Cambridge University Press, 2002.
- [6] P A Calogero, P R Gorda, A Pickering. *Multicomponent equations associated to non-isospectral scattering problems*, Inverse Probl, 1997, 13(6): 1463-1476.
- [7] A Sakhnovich. *Nonisospectral integrable nonlinear equations with external potentials and their GBDT solutions*, J Phys A: Math Theor, 2008, 155: 155204.
- [8] M J Ablowitz, H Segur. *Solitons and the Inverse Scattering Transform*, Philadelphia, PA: SIAM, 1981.

- [9] A C Newell. *Solitons in Mathematics and Physics*, Philadelphia, PA: SIAM, 1985.
- [10] F Magri. *Nonlinear Evolution Equations and Dynamical Systems*, In: Lecture Notes in Physics, vol 120, Berlin: Springer, 1980.
- [11] G Z Tu. *The trace identity, a powerful tool for constructing the Hamiltonian structure of integrable systems*, J Math Phys, 1989, 30(2): 330-338.
- [12] Y Zhang, W Rui. *A few continuous and discrete dynamical systems*, Rep Math Phys, 2016, 78(1): 19-32.
- [13] Y Zhang, H Tam. *Discussion on integrable properties for higher-dimensional variable-coefficient nonlinear partial differential equations*, J Math Phys, 2013, 54(1): 013516.
- [14] Y Zhang, E Fan, H Tam. *A few expanding Lie algebras of the Lie algebra A_1 and applications*, Phys Lett A, 2006, 359(5): 471-480.
- [15] Y F Zhang, J Liu. *Induced Lie algebras of a six-dimensional matrix Lie algebra*, Commun Theor Phys, 2008, 50(2): 289.
- [16] Z Qiao. *New hierarchies of isospectral and non-isospectral integrable NLEEs derived from the Harry-Dym spectral problem*, Physica A, 1998, 252(34): 377-387.
- [17] Y S Li, G C Zhu. *New set of symmetries of the integrable equations, Lie algebras and non-isospectral evolution equations. II. AKNS system*, J Phys A: Math Gen, 1986, 19(8): 3713-3725.
- [18] X X Xu. *An integrable coupling hierarchy of the Mkdv-integrable systems, its Hamiltonian structure and corresponding nonisospectral integrable hierarchy*, Appl Math Comput, 2010, 216(1): 344-353.
- [19] Y Zhang, H Zhang, Q Yan. *Integrable couplings of Botie-Pempinelli-Tu (BPT) hierarchy*, Phys Lett A, 2002, 299(5-6): 543-548.
- [20] W X Ma. *An approach for constructing non-isospectral hierarchies of evolution equations*, J Phys A: Math Gen, 1992, 25: L719-L726.
- [21] W X Ma. *A simple scheme for generating nonisospectral flows from the zero curvature representation*, Phys Lett A, 1993, 179: 179-185.
- [22] W X Ma. *K symmetries and τ symmetries of evolution equations and their Lie algebras*, J Phys A: Math Gen, 1990, 23: 2707-2716.
- [23] Z Qiao. *Generation of soliton hierarchy and general structure of its commutator representations*, Acta Math Appl Sin, 1995, 18(2): 287-301.(in Chinese)
- [24] Z Qiao. *Algebraic structure of the operator related to stationary systems*, Phys Lett A, 1995, 206(5-6): 347-358.
- [25] M Boiti, F Pempinelli, G Z Tu. *Canonical structure of soliton equations via isospectral eigenvalue problems*, Nuovo Cim B, 1984, 79: 231-265.
- [26] M Boiti, G Z Tu. *A simple approach to the Hamiltonian structure of soliton equations*, Nuovo Cim B, 1983, 75: 145.
- [27] Y Zhang, J Mei, H Guan. *A method for generating isospectral and nonisospectral hierarchies of equations as well as symmetries*, J Geom Phys, 2020, 147: 103538.
- [28] Y Zhang, H Tam. *Three kinds of coupling integrable couplings of the Korteweg-de Vries hierarchy of evolution equations*, J Math Phys, 2010, 51(4): 043510.
- [29] W X Ma, X X Xu, Y Zhang. *Semi-direct sums of Lie algebras and discrete integrable couplings*, J Math Phys, 2006, 47: 053501.

- [30] W X Ma. *Integrable couplings and matrix loop algebras*, AIP Conf Proc, 2013, 1562: 105-122.
- [31] W X Ma. *The algebraic structure of zero curvature representations and application to coupled KdV systems*, J Phys A: Math Gen, 1993, 26(11): 2573-2582.
- [32] W X Ma. *The algebraic structures of isospectral Lax operators and applications to integrable equations*, J Phys A: Math Gen, 1992, 25(20): 5329-5343.
- [33] H Wang, C Li. *Affine Weyl group symmetries of Frobenius Painlevé equations*, Math Meth Appl Sci, 2020, 436: 3238-3252.
- [34] P G Estévez, C Savdón. *Miura-reciprocal transformations for non-isospectral Camassa-Holm hierarchies in 2+1 dimensions*, J Nonlinear Math Phys, 2013, 20(4): 552-564.
- [35] X H Zhao, B Tian, H M Li, et al. *Solitons, periodic waves, breathers and integrability for a nonisospectral and variable-coefficient fifth-order Korteweg-de Vries equation in fluids*, Appl Math Lett, 2017, 65: 48-55.
- [36] F Yu. *A novel non-isospectral hierarchy and soliton wave dynamics for a parity-time-symmetric nonlocal vector nonlinear Gross-Pitaevskii equations*, Commun Nonlinear Sci Numer Simul, 2019, 78: 104852.
- [37] X Y Gao, Y J Guo, W R Shan. *Optical waves/modes in a multicomponent inhomogeneous optical fiber via a three-coupled variable-coefficient nonlinear Schrödinger system*, Appl Math Lett, 2021, 120: 107161.
- [38] X Y Gao, Y J Guo, W R Shan. *Bilinear forms through the binary Bell polynomials, N solitons and Bäcklund transformations of the Boussinesq-Burgers system for the shallow water waves in a lake or near an ocean beach*, Commun Theor Phys, 2020, 72: 095002.
- [39] X Y Gao, Y J Guo, W R Shan. *Looking at an open sea via a generalized (2+1)-dimensional dispersive long-wave system for the shallow water: scaling transformations, hetero-Bäcklund transformations, bilinear forms and N solitons*, Eur Phys J Plus, 2021, 136: 893.
- [40] X Y Gao, Y J Guo, W R Shan. *Symbolic computation on a (2+1)-dimensional generalized variable-coefficient Boiti-Leon-Pempinelli system for the water waves*, Chaos Solitons Fract, 2021, 150: 111066.
- [41] M Wang, B Tian, C C Hu, et al. *Generalized Darboux transformation, solitonic interactions and bound states for a coupled fourth-order nonlinear Schrödinger system in a birefringent optical fiber*, Appl Math Lett, 2021, 119: 106936.
- [42] Y Shen, B Tian. *Bilinear auto-Bäcklund transformations and soliton solutions of a (3+1)-dimensional generalized nonlinear evolution equation for the shallow water waves*, Appl Math Lett, 2021, 122: 107301.
- [43] X T Gao, B Tian, Y Shen, et al. *Comment on “Shallow water in an open sea or a wide channel: Auto- and non-auto-Bäcklund transformations with solitons for a generalized (2+1)-dimensional dispersive long-wave system”*, Chaos Solitons Fract, 2021, 151: 111222.
- [44] D Y Yang, B Tian, Q Y Qu, et al. *Lax pair, conservation laws, Darboux transformation and localized waves of a variable-coefficient coupled Hirota system in an inhomogeneous optical fiber*, Chaos Solitons Fract, 2021, 150: 110487.

¹School of Science, Jimei University, Xiamen 361021, China.

²School of Mathematics, China University of Mining and Technology, Xuzhou 221116, China.

E-mail: zhangyfcumt@163.com