

# On the monotonicity of limit wave speed to a perturbed gKdV equation

WEN Zhen-shu\*      SHI Tian-yu

**Abstract.** This paper deals with the monotonicity of limit wave speed  $c_0(h)$  to a perturbed gKdV equation. We show the decrease of  $c_0(h)$  by combining the analytic method and the numerical technique. Our results solve a special case of the open question presented by Yan et al., and the method potentially provides a way to study the monotonicity of  $c_0(h)$  for general  $m \in \mathbb{N}^+$ .

## §1 Introduction

Many shallow water wave models have been found in the field of water wave dynamics, which play an important role in describing natural phenomena [1–4]. In recent years, there exist more and more works concerning the perturbed models and their solutions due to their realistic nature [5–13], among which, the perturbed KdV equation and its variants have gained considerate concern since they are significant in physics and waves contexts.

In 2014, Yan et al. [14] showed the persistence of solitary waves and periodic waves to the perturbed generalized KdV (gKdV) equation

$$v_t + v^m v_x + v_{xxx} + \sigma(v_{xx} + v_{xxxx}) = 0, \quad (1)$$

for general positive integer  $m$  and sufficiently small  $\sigma > 0$ . When  $m = 1$ , we have the perturbed KdV equation

$$v_t + vv_x + v_{xxx} + \sigma(v_{xx} + v_{xxxx}) = 0, \quad (2)$$

which was investigated by Owaga [15]. Owaga [15] derived the existence of the above two types of waves to (2), and showed that  $c_0(h)$  is decreasing.

Naturally, one may wonder how about the monotonicity of  $c_0(h)$  to Eq.(1) with general  $m \in \mathbb{N}^+$ , which was presented as an open question in [14]. In fact, the monotonicity of  $c_0(h)$  for  $m = 1, 2, 3, 4$  had been shown [15–17]. However, as suggested in [17], for  $m \geq 5$ , it is difficult

---

Received: 2022-01-26.      Revised: 2022-04-25.

MR Subject Classification: 65L11, 35C07, 33F05.

Keywords: the perturbed gKdV equation with  $m = 5$ , traveling waves, limit wave speed, monotonicity.

Digital Object Identifier(DOI): <https://doi.org/10.1007/s11766-025-4675-1>.

Supported by the National Natural Science Foundation of China(12071162), the Natural Science Foundation of Fujian Province(2021J01302), and the Fundamental Research Funds for the Central Universities(ZQN-802).

\*Corresponding author.

to show  $(J'_m/J'_0)' > 0$  analytically ( $J_m$  is the Abelian integral defined later), which is a crucial result for later analysis, due to the complexity of the Riccati equations. To find a way to address the monotonicity of  $c_0(h)$  for odd  $m$ , in this paper, we first focus on a special case when  $m = 5$ .

By the transformation  $\zeta = x - ct$  with the wave speed  $c > 0$ , Eq.(1) with  $m = 5$  can be converted into

$$-cv'(\zeta) + v^5(\zeta)v'(\zeta) + v'''(\zeta) + \sigma(v''(\zeta) + v''''(\zeta)) = 0, \quad (3)$$

and it follows that

$$-cv(\zeta) + \frac{1}{6}v^6(\zeta) + v''(\zeta) + \sigma(v'(\zeta) + v'''(\zeta)) = 0, \quad (4)$$

by integrating.

Introducing  $\eta = \sqrt{c}\zeta$ ,  $v = \sqrt[5]{c}X$ , we obtain

$$-X(\eta) + \frac{1}{6}X^6(\eta) + X''(\eta) + \sigma\left(\frac{1}{\sqrt{c}}X'(\eta) + \sqrt{c}X'''(\eta)\right) = 0. \quad (5)$$

If we set  $\sigma = 0$ , then (5) becomes the unperturbed system

$$X'' + \frac{1}{6}X^6 - X = 0, \quad (6)$$

which can be rewritten as a planar system

$$\begin{cases} \frac{dX}{d\eta} = Y, \\ \frac{dY}{d\eta} = X - \frac{1}{6}X^6, \end{cases} \quad (7)$$

whose first integral is

$$H(X, Y) = -Y^2 + X^2 - \frac{1}{21}X^7 = h. \quad (8)$$

Note that system (7) has one saddle  $(0, 0)$  and one center  $(\sqrt[5]{6}, 0)$ , and its phase portrait is shown in Figure 1. Besides, the traveling waves of the system (7) can be parameterized through  $h$ . Exploiting this parametrization, we give the monotonicity of  $c_0(h)$  in Theorem 1.

**Theorem 1.** For  $h \in [0, \frac{5\sqrt[5]{36}}{7})$ , the perturbed gKdV equation (1) with  $m = 5$  and sufficiently small  $\sigma > 0$  has a traveling wave  $v = \sqrt[5]{c}X(\sigma, h, c, \eta)$ , where  $X(\sigma, h, c, \eta)$  is the solution of (5), and  $c = c(\sigma, h)$  depends on  $\sigma$  and  $h$  smoothly, with the limit  $c_0(h)$  as  $\sigma \rightarrow 0$ . Furthermore,  $c_0(h)$  is a smooth decreasing function for  $h \in [0, \frac{5\sqrt[5]{36}}{7})$ .

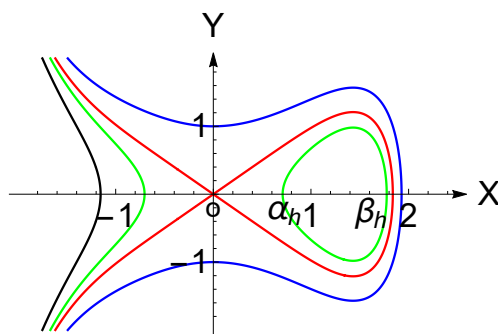


Figure 1. The phase portrait of system (7).

## §2 The theoretic derivations of the monotonicity of $c_0(h)$

In this section, we focus on theoretic derivations of the monotonicity of  $c_0(h)$  by applying the Abelian integral theory and numerical technique, see [14] for other results. Suppose that  $X(\eta)$  is the solution of the system (7), and we define  $Q$  and  $R$  as

$$Q = \frac{1}{2} \int (X'')^2 d\eta, \quad R = \frac{1}{2} \int (X')^2 d\eta,$$

where the integrals are performed along the orbits of system (7). Then from [14], it is known that  $c_0(h)$  can be expressed as

$$c_0(h) = \frac{1}{Z(h)},$$

where

$$Z(h) = \frac{Q}{R}.$$

Assume that  $\alpha(h)$  and  $\beta(h)$  ( $0 \leq \alpha(h) < \beta(h)$ ) are the two roots of  $X^2 - \frac{1}{21}X^7 = h$ , where  $0 \leq h < \frac{5\sqrt[5]{36}}{7}$ . Therefore,  $Q$  and  $R$  can be rewritten as

$$Q = \int_{\alpha(h)}^{\beta(h)} \frac{(X - \frac{1}{6}X^6)^2}{E(X)} dX, \quad R = \int_{\alpha(h)}^{\beta(h)} E(X) dX, \quad (9)$$

with  $E(X) = \sqrt{X^2 - \frac{1}{21}X^7 - h}$ .

Introducing the Abelian integrals:

$$J_n(h) = \int_{\alpha(h)}^{\beta(h)} X^n E(X) dX, \quad n = 0, 1, 2, \dots, \quad (10)$$

which satisfy

$$\int_{\alpha(h)}^{\beta(h)} \frac{X^n}{E(X)} dX = -2J'_n(h), \quad (11)$$

by direct computation.

To prove that  $c_0(h)$  is decreasing, we only need to show that  $Z(h)$  is increasing. Here we need Lemmas 1 and 2 from [14].

**Lemma 1.** [14]. We have  $J = A(h)J'$ , where  $J = (J_0, J_1, J_2, J_3, J_4, J_5)^T$  and

$$A(h) = \begin{pmatrix} \frac{14}{9}h & 0 & -\frac{10}{9} & 0 & 0 & 0 \\ 0 & \frac{14}{11}h & 0 & -\frac{10}{11} & 0 & 0 \\ 0 & 0 & \frac{14}{13}h & 0 & -\frac{10}{13} & 0 \\ 0 & 0 & 0 & \frac{14}{15}h & 0 & -\frac{10}{15} \\ 0 & -\frac{60}{17} & 0 & 0 & \frac{14}{17}h & 0 \\ \frac{140}{57}h & 0 & -\frac{280}{57} & 0 & 0 & \frac{14}{19}h \end{pmatrix}.$$

**Lemma 2.** [14]. We have  $Q = J_5 - J_0, R = J_0$  and  $Z(h) = \frac{J_5}{J_0} - 1$ .

From Lemma 1, we can derive the inverse of  $A$ ,

$$A^{-1} = \frac{1}{\Delta} \times \begin{pmatrix} 9((14h)^5 - 56 \cdot 10^5) & 66 \cdot 10^3(14h)^2 & 130(14h)^4 & 9 \cdot 10^5 \cdot 14h & 17 \cdot 10^2(14h)^3 & 114 \cdot 10^5 \\ -3 \cdot 10^3(14h)^3 & 11(14h)^5 & 78 \cdot 10^3(14h)^2 & 150(14h)^4 & 102 \cdot 10^4 \cdot 14h & 19 \cdot 10^2(14h)^3 \\ -18 \cdot 10^5 \cdot 14h & 66 \cdot 10^2(14h)^3 & 13(14h)^5 & 9 \cdot 10^4(14h)^2 & 170(14h)^4 & 114 \cdot 10^4 \cdot 14h \\ -3 \cdot 10^2(14h)^4 & 396 \cdot 10^4 14h & 78 \cdot 10^2(14h)^3 & 15(14h)^5 & 102 \cdot 10^3(14h)^2 & 190(14h)^4 \\ -18 \cdot 10^4(14h)^2 & 660(14h)^4 & 468 \cdot 10^4 14h & 9 \cdot 10^3(14h)^3 & 17(14h)^5 & 114 \cdot 10^3(14h)^2 \\ -30(14h)^5 & 396 \cdot 10^3(14h)^2 & 780(14h)^4 & 54 \cdot 10^5 \cdot 14h & 102 \cdot 10^2(14h)^3 & 19(14h)^5 \end{pmatrix}, \quad (12)$$

with  $\Delta = 14h((14h)^5 - 36 \times 10^5)$ .

With these preparations, now we will prove that  $Z(h)$  is increasing. Exploiting Lemma 1 and (12), we derive the following Ricatti equation

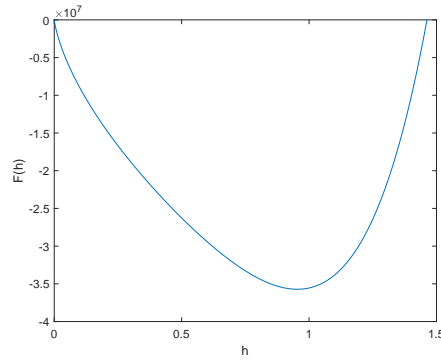
$$\begin{aligned} Z'(h) &= \left( \frac{J_5}{J_0} \right)' = \frac{J_5'}{J_0} - \frac{J_0' J_5}{(J_0)^2} \\ &= \frac{1}{\Delta J_0} (-30(14h)^5 J_0 + 396 \cdot 10^3(14h)^2 J_1 + 780(14h)^4 J_2 + 54 \cdot 10^5 \cdot 14h J_3 + 102 \cdot 10^2(14h)^3 J_4 \\ &\quad + 19(14h)^5 J_5) - \frac{J_5}{\Delta(J_0)^2} (9((14h)^5 - 56 \cdot 10^5) J_0 + 66 \cdot 10^3(14h)^2 J_1 + 130(14h)^4 J_2 \\ &\quad + 9 \cdot 10^5 \cdot 14h J_3 + 17 \cdot 10^2(14h)^3 J_4 + 114 \cdot 10^5 J_5) \\ &= \frac{1}{\Delta} \left( -30(14h)^5 + 396 \cdot 10^3(14h)^2 \frac{J_1}{J_0} + 780(14h)^4 \frac{J_2}{J_0} + 54 \cdot 10^5 \cdot 14h \frac{J_3}{J_0} + 102 \cdot 10^2(14h)^3 \frac{J_4}{J_0} \right. \\ &\quad + (10(14h)^5 + 504 \cdot 10^5) \frac{J_5}{J_0} - 66 \cdot 10^3(14h)^2 \frac{J_1}{J_0} \frac{J_5}{J_0} + 130(14h)^4 \frac{J_2}{J_0} \frac{J_5}{J_0} - 9 \cdot 10^5 \cdot 14h \frac{J_3}{J_0} \frac{J_5}{J_0} \\ &\quad \left. - 17 \cdot 10^2(14h)^3 \frac{J_4}{J_0} \frac{J_5}{J_0} - 114 \cdot 10^5 \left( \frac{J_5}{J_0} \right)^2 \right) \\ &= \frac{1}{\Delta} F(h), \end{aligned} \quad (13)$$

where

$$\begin{aligned} F(h) &= -114 \cdot 10^5 \left( \frac{J_5}{J_0} \right)^2 + (10(14h)^5 + 504 \cdot 10^5) \frac{J_5}{J_0} - 30(14h)^5 \\ &\quad + \frac{140h}{J_0} \left( 6 - \frac{J_5}{J_0} \right) (6600 \cdot 14h J_1 + 13(14h)^3 J_2 + 9 \cdot 10^4 J_3 + 170(14h)^2 J_4). \end{aligned} \quad (14)$$

Since  $\Delta$  is negative, if we can prove  $F(h) < 0$ , then  $Z'(h) > 0$ , and it follows that  $c_0(h)$  is decreasing. However, it is not easy to directly show that  $F(h) < 0$  analytically. Therefore, here we develop a numerical technique to show it. Specifically, let  $h$  change from 0.0001 to 1.4626 (note that  $\frac{5\sqrt[5]{36}}{7} \approx 1.4626$ ) with step 0.0001, then we can numerically evaluate  $\alpha(h), \beta(h)$  and  $J_i, i = 0, 1, \dots, 5$  through (10), and the values of  $F(h)$  follow. The profile of  $F(h)$  with respect to  $h$  is shown in Figure 2 and the maximum of  $F(h)$  is  $-4.2866 \times 10^3$  at  $h = 1.4626$ , which indicates that  $F(h) < 0$ , for  $0 < h < \frac{5\sqrt[5]{36}}{7}$ .

**Remark 1.** As implied in [17], when  $m \geq 5$ , it is very difficult to prove the monotonicity of  $c_0(h)$  analytically. Therefore, we turn to the numerical technique and solve the open question

Figure 2. The profile of  $F(h)$  with respect to  $h$ .

for the special case when  $m = 5$ .

### §3 Conclusions

In this paper, we develop the numerical technique to derive the monotonicity of  $c_0(h)$  for a special perturbed gKdV equation (1) when  $m = 5$ , which partially solves the open question presented by [14]. More importantly and interestingly, it potentially provides a way to answer the open question for general  $m \in \mathbb{N}^+$  completely, which is still under consideration. The main difficulties lie in two aspects, the complexity of  $A^{-1}$  and  $Z'(h)$ , and the algorithm and program.

### Declarations

**Conflict of interest** The authors declare no conflict of interest.

### References

- [1] Z Wen. *Existence and dynamics of bounded traveling wave solutions to Getmanou equation*, Communications in Theoretical Physics, 2018, 70: 672-676.
- [2] Z Wen. *The generalized bifurcation method for deriving exact solutions of nonlinear space-time fractional partial differential equations*, Applied Mathematics and Computation, 2021, 366: 124735.
- [3] Z Wen. *Qualitative study of effects of vorticity on traveling wave solutions to the two-component Zakharov-Itō system*, Applicable Analysis, 2021, 100(11): 2334-2346.
- [4] Z Wen, G Chen, J Li. *Pseudo-peakon, periodic peakons and compactons on a shallow water model with the Coriolis effect*, International Journal of Bifurcation and Chaos, 2021, 31(8): 2150144.

- [5] A Chen, L Guo, X Deng. *Existence of solitary waves and periodic waves for a perturbed generalized BBM equation*, Journal of Differential Equations, 2016, 261(10): 5324-5349.
- [6] Z Wen. *On existence of kink and antikink wave solutions of singularly perturbed Gardner equation*, Mathematical Methods in the Applied Sciences, 2020, 43(7): 4422-4427.
- [7] Z Wen, L Zhang, M Zhang. *Dynamics of classical Poisson-Nernst-Planck systems with multiple cations and boundary layers*, Journal of Dynamics and Differential Equations, 2021, 33: 211-234.
- [8] J Ge, Z Du. *The solitary wave solutions of the nonlinear perturbed shallow water wave model*, Applied Mathematics Letters, 2020, 103: 106202.
- [9] L Guo, Y Zhao. *Existence of periodic waves for a perturbed quintic BBM equation*, Discrete and Continuous Dynamical Systems, 2020, 40(8): 4689-4703.
- [10] Z Wen, H Li, Y Fu. *Abundant explicit periodic wave solutions and their limit forms to space-time fractional Drinfel'd-Sokolov-Wilson equation*, Mathematical Methods in the Applied Sciences, 2021, 44: 6406-6421.
- [11] H Zhang, Y Xia, P R N'gbo. *Global existence and uniqueness of a periodic wave solution of the generalized Burgers-Fisher equation*, Applied Mathematics Letters, 2021, 121: 107353.
- [12] Z Wen, P W Bates, M Zhang. *Effects on I-V relations from small permanent charge and channel geometry via classical Poisson-Nernst-Planck equations with multiple cations*, Nonlinearity, 2021, 34: 4464-4502.
- [13] P W Bates, Z Wen, M Zhang. *Small permanent charge effects on individual fluxes via classical Poisson-Nernst-Planck systems with multiple cations*, Journal of Nonlinear Science, 2021, 33: 55.
- [14] W Yan, Z Liu, Y Liang. *Existence of solitary waves and periodic waves to a perturbed generalized KdV equation*, Mathematical Modelling and Analysis, 2014, 19(4): 537-555.
- [15] T Ogawa. *Travelling wave solutions to a perturbed Korteweg-de Vries equation*, Hiroshima Mathematical Journal, 1994, 24(2): 401-422.
- [16] A Chen, C Zhang, W Huang. *Limit speed of traveling wave solutions for the perturbed generalized KdV equation*, Discrete and Continuous Dynamical Systems-S, 2022, 16: 379-402.
- [17] A Chen, C Zhang, W Huang. *Monotonicity of limit wave speed of traveling wave solutions for a perturbed generalized KdV equation*, Applied Mathematics Letters, 2021, 121: 107381.

School of Mathematical Sciences, Huaqiao University, Quanzhou 362021, China.

Email: wenzhenshu@hqu.edu.cn